

Tutorial sheet 13

Discussion topic: What are the fundamental equations of the dynamics of a relativistic fluid? What is the relation between the energy-momentum tensor of a perfect relativistic fluid and its internal energy, pressure, and four-velocity? How is the latter defined?

Hint: If the covariant derivatives d_μ in the following exercises upset you, choose Minkowski coordinates, in which $d_\mu = \partial_\mu$.

34. Quantum number conservation

Consider a 4-current with components $N^\mu(x)$ obeying the continuity equation $d_\mu N^\mu(x) = 0$. Show that the quantity $\mathcal{N} = \int N^0(x) d^3\vec{r}/c$ is a Lorentz scalar, by convincing yourself first that \mathcal{N} can be rewritten in the form

$$\mathcal{N} = \frac{1}{c} \int_{x^0=\text{const.}} N^\mu(x) d^3\sigma_\mu, \tag{1}$$

where $d^3\sigma_\mu = \frac{1}{6} \epsilon_{\mu\nu\rho\sigma} d^3\mathcal{V}^{\nu\rho\sigma}$ is a 4-vector, with $d^3\mathcal{V}^{\nu\rho\sigma}$ the antisymmetric 4-tensor defined by

$$d^3\mathcal{V}^{012} = dx^0 dx^1 dx^2, \quad d^3\mathcal{V}^{021} = -dx^0 dx^2 dx^1, \quad \text{etc.}$$

and $\epsilon_{\mu\nu\rho\sigma}$ the totally antisymmetric Levi-Civita tensor with the convention $\epsilon_{0123} = +1$, such that $d^3\mathcal{V}^{\nu\rho\sigma}$ represents the 3-dimensional hypersurface element in Minkowski space.

35. Energy-momentum tensor

Let \mathcal{R} denote a fixed reference frame. Consider a perfect fluid whose local rest frame at a point x moves with velocity \vec{v} with respect to \mathcal{R} . Show with the help of a Lorentz transformation that the Minkowski components of the energy-momentum tensor of the fluid at x are given to order $\mathcal{O}(|\vec{v}|/c)$ by

$$T^{\mu\nu} = \begin{pmatrix} \epsilon & (\epsilon + \mathcal{P})\frac{v^1}{c} & (\epsilon + \mathcal{P})\frac{v^2}{c} & (\epsilon + \mathcal{P})\frac{v^3}{c} \\ (\epsilon + \mathcal{P})\frac{v^1}{c} & \mathcal{P} & 0 & 0 \\ (\epsilon + \mathcal{P})\frac{v^2}{c} & 0 & \mathcal{P} & 0 \\ (\epsilon + \mathcal{P})\frac{v^3}{c} & 0 & 0 & \mathcal{P} \end{pmatrix},$$

where for the sake of brevity the x -dependence of the various fields is omitted. Check the compatibility of this result with the general formula for $T^{\mu\nu}$ given in the lecture.

36. Equations of motion of a perfect relativistic fluid

In this exercise, we set $c = 1$ and drop the x variable for the sake of brevity. Remember that the metric tensor has signature $(-, +, +, +)$.

i. Check that the tensor with components $\Delta^{\mu\nu} \equiv g^{\mu\nu} + u^\mu u^\nu$ defines a projector on the subspace orthogonal to the 4-velocity.

Denoting by d_μ the components of the (covariant) 4-gradient, we define $\nabla^\nu \equiv \Delta^{\mu\nu} d_\mu$. Can you see the rationale behind this notation?

ii. Show that the energy-momentum conservation equation for a perfect fluid is equivalent to the two equations

$$u^\mu d_\mu \epsilon + (\epsilon + \mathcal{P}) d_\mu u^\mu = 0 \quad \text{and} \quad (\epsilon + \mathcal{P}) u^\mu d_\mu u^\nu + \nabla^\nu \mathcal{P} = 0. \tag{2}$$

Which known equation does the second one evoke?

37. A family of solutions of the dynamical equations for perfect relativistic fluids

Let $\{x^\mu\}$ denote Minkowski coordinates and $\tau^2 \equiv -x^\mu x_\mu$, where the “mostly plus” metric is used. Show that the following four-velocity, pressure and charge density constitute a solution of the equations describing the motion of a perfect relativistic fluid with equation of state $\mathcal{P} = K\varepsilon$ and a single conserved charge:

$$u^\mu(x) = \frac{x^\mu}{\tau} \quad , \quad \mathcal{P}(x) = \mathcal{P}_0 \left(\frac{\tau_0}{\tau} \right)^{3(1+K)} \quad , \quad n(x) = n_0 \left(\frac{\tau_0}{\tau} \right)^3 \mathcal{N}(\sigma(x)), \quad (3)$$

with τ_0 , \mathcal{P}_0 , n_0 arbitrary constants and \mathcal{N} an arbitrary function of a single argument, while σ is a function of spacetime coordinates with vanishing comoving derivative: $u^\mu \partial_\mu \sigma(x) = 0$.