

Curvilinear coordinates


Curvilinear coordinates

- Changing coordinates
- Covariant differentiation of vector and tensor fields
- Application (tomorrow): Milne coordinates

Recipes, without proof nor attempt at rigor!

Motivation: an issue with polar coordinates

Consider the (2D) plane with

● Cartesian coordinates $(x,y) \stackrel{\text{def}}{=} (x^1,x^2)$ with associated basis vectors $(\vec{e}_x,\vec{e}_y) \stackrel{\text{def}}{=} (\vec{e}_1,\vec{e}_2)$ – note the at first bizarre primed indices as well as their positions;  the reference!

● the usual polar coordinates $(r,\theta) \stackrel{\text{def}}{=} (x^1,x^2)$ with the associated basis vectors $(\vec{e}_r,\vec{e}_\theta) \stackrel{\text{def}}{=} (\vec{e}_1,\vec{e}_2)$ – here with unprimed indices.

The position \vec{R} of a point P of the plane may be denoted

$$\vec{R} = x\vec{e}_x + y\vec{e}_y \equiv \sum x^i\vec{e}_i \equiv x^i\vec{e}_i$$

But you know that you would NOT write $\vec{R} = r\vec{e}_r + \theta\vec{e}_\theta \equiv x^i\vec{e}_i$. 

Motivation: an issue with polar coordinates

If you look at two neighboring points P and $P + dP$ and call $d\vec{R}$ their separation, you will also write (Cartesian coordinates)

$$d\vec{R} = dx\vec{e}_x + dy\vec{e}_y \equiv \sum dx^i\vec{e}_i \equiv dx^i\vec{e}_i'$$

And perhaps are you then ready to write (polar coordinates)

$$d\vec{R} = dr\vec{e}_r + d\theta\vec{e}_\theta$$

so that at least one formula:

$$d\vec{R} = dx^i\vec{e}_i$$

(1)

is valid in every coordinate system.

Fine? Let us see where this is leading to...

Separation vector between neighboring points in arbitrary coordinates

It is tempting to write $d\vec{R} = \frac{\partial \vec{R}}{\partial x^i} dx^i$.

An issue is that we know how to express \vec{R} as a function of the Cartesian coordinates $\{x^{i'}\}$, but in general not in terms of arbitrary coordinates $\{x^i\}$.

We generally know how to express the $\{x^i\}$ in terms of the $\{x^{i'}\}$ and reciprocally: functions $x^i(\{x^{i'}\})$ and $x^{i'}(\{x^i\})$.

For instance: $r(x, y) = \sqrt{x^2 + y^2}$, $\theta(x, y) = \arctan(y/x)$

and $x(r, \theta) = r \cos \theta$, $y(r, \theta) = r \sin \theta$

Separation vector between neighboring points in arbitrary coordinates

It is tempting to write $d\vec{R} = \frac{\partial \vec{R}}{\partial x^i} dx^i$. (2)

An issue is that we know how to express \vec{R} as a function of the Cartesian coordinates $\{x^{i'}\}$, but in general not in terms of arbitrary coordinates $\{x^i\}$.

We generally know how to express the $\{x^i\}$ in terms of the $\{x^{i'}\}$ and reciprocally: functions $x^i(\{x^{i'}\})$ and $x^{i'}(\{x^i\})$.

👉 We may write $\vec{R}(\{x^i\}) = x(\{x^i\})\vec{e}_x + y(\{x^i\})\vec{e}_y = x^{i'}(\{x^i\})\vec{e}_{i'}$
and now differentiate w.r.t. x^i :

$$\frac{\partial \vec{R}}{\partial x^i} = \frac{\partial x^{i'}}{\partial x^i} \vec{e}_{i'} \quad \Rightarrow \quad d\vec{R} = \frac{\partial x^{i'}}{\partial x^i} \vec{e}_{i'} dx^i \quad (3)$$

Basis vectors of an arbitrary coordinate system

Identifying Eqs. (1) and (3), we find the “proper” basis vectors $\{\vec{e}_i\}$:

$$\left. \begin{aligned} d\vec{R} &= dx^i \vec{e}_i \\ d\vec{R} &= \frac{\partial x^{i'}}{\partial x^i} \vec{e}_{i'} dx^i \end{aligned} \right\} \vec{e}_i = \frac{\partial x^{i'}}{\partial x^i} \vec{e}_{i'} \quad (4)$$

For the polar coordinates: $x(r, \theta) = r \cos \theta$, $y(r, \theta) = r \sin \theta$

$$\frac{\partial x}{\partial r} = \cos \theta , \quad \frac{\partial y}{\partial r} = \sin \theta \quad \Rightarrow \quad \vec{e}_r = \cos \theta \vec{e}_x + \sin \theta \vec{e}_y$$

$$\frac{\partial x}{\partial \theta} = -r \sin \theta , \quad \frac{\partial y}{\partial \theta} = r \cos \theta \quad \Rightarrow \quad \vec{e}_\theta = -r \sin \theta \vec{e}_x + r \cos \theta \vec{e}_y$$

Well \vec{e}_r is fine and I knew it, but... this \vec{e}_θ is not normalized to 1! 🤔

Motivation: a problem with polar coordinates

For the polar coordinates: $x(r, \theta) = r \cos \theta$, $y(r, \theta) = r \sin \theta$

$$\vec{e}_r = \cos \theta \vec{e}_x + \sin \theta \vec{e}_y \qquad \vec{e}_\theta = -r \sin \theta \vec{e}_x + r \cos \theta \vec{e}_y$$

Well \vec{e}_r is fine and I knew it, but... this \vec{e}_θ is not normalized to 1! 🤔

Indeed! But who said it should? 🧐

Now \vec{e}_θ is well-behaved in the limit $r \rightarrow 0$, when θ is not defined: isn't this a nice feature, worth discarding the normalization? 😇

(the perplexed students):

OK, we'll live with it...

Metric tensor

Start from eq.(1): $d\vec{R} = dx^i \vec{e}_i$ and square it:

$$d\vec{R}^2 = (dx^i \vec{e}_i) \cdot (dx^j \vec{e}_j) = (\vec{e}_i \cdot \vec{e}_j) dx^i dx^j$$

Define $g_{ij} \equiv \vec{e}_i \cdot \vec{e}_j$ and replace: $d\vec{R}^2 = g_{ij} dx^i dx^j$

● Cartesian coordinates: $g_{i'j'} = \delta_{i'j'} \Leftrightarrow d\vec{R}^2 = (dx)^2 + (dy)^2$

● Polar coordinates:* $g_{11} = g_{rr} = 1$, $g_{22} = g_{\theta\theta} = r^2$, $g_{r\theta} = g_{\theta r} = 0$
 $\Leftrightarrow d\vec{R}^2 = (dr)^2 + r^2(d\theta)^2$

looks familiar...

g_{ij} is called **metric tensor**.

(More precisely, these are its components, but let's not be picky today).

* Using the expressions of \vec{e}_r , \vec{e}_θ in terms of \vec{e}_x , \vec{e}_y , since we know the scalar product in Cartesian coordinates.

Metric tensor and its inverse

Define g^{ij} such that its “product” (more accurately: contraction) with the metric tensor equals the identity:

$$g^{ij}g_{jk} = \delta_k^i$$

g^{ij} is called the inverse metric tensor.

Cartesian coordinates: $g^{i'j'} = \delta^{i'j'}$; polar coordinates: $g^{rr} = 1$, $g^{\theta\theta} = 1/r^2$.

g_{ij} and g^{ij} can be used to lower / raise indices:

- No impact on Cartesian components
- In polar coordinates: no impact on r -components, θ -components vary.

👉 What does it mean? Not much today...

Next step: generalize!

Starting from coordinates $\{x^{i'}\}$ with orthogonal* and normalized* basis vectors $\{\vec{e}_{i'}\}$, one goes to alternative coordinates $\{x^i\}$:

- with the basis vectors $\vec{e}_i = \frac{\partial x^{i'}}{\partial x^i} \vec{e}_{i'}$ such that $d\vec{R} = dx^i \vec{e}_i$ holds
- and the metric tensor $g_{ij} \equiv \vec{e}_i \cdot \vec{e}_j$ such that $d\vec{R}^2 = g_{ij} dx^i dx^j$.

We may do that:

- ... in 2-dimensional Euclidean space: $g_{i'j'} = \text{diag}(+1, +1)$
- ... in 3-dimensional Euclidean space: $g_{i'j'} = \text{diag}(+1, +1, +1)$
- ... in 4-dimensional Minkowski space: $g_{i'j'} = \text{diag}(-1, +1, +1, +1)$

* Using the appropriate (pseudo)scalar product encoded in $g_{i'j'}$

Second motivation: another issue with polar coordinates

How should one differentiate?

Consider a (smooth enough) function f .

Its derivatives are supposed to measure the variations of the function in the various respective directions:

$$f(x + \delta x, y) \simeq f(x, y) + \frac{\partial f}{\partial x} \delta x$$

and so on...

Second motivation: another issue with polar coordinates

How should one differentiate?

Consider a (smooth enough) function $f(r, \theta)$.

Its derivatives are supposed to measure the variations of the function in the various respective directions:

$$f(r, \theta + \delta\theta) \simeq f(r, \theta) + \frac{\partial f}{\partial \theta} \delta\theta$$

and so on...

Second motivation: another issue with polar coordinates

For instance: $\vec{f}(r, \theta) = \cos \theta \vec{e}_r - \frac{\sin \theta}{r} \vec{e}_\theta$

with (slide 6) $\vec{e}_r = \cos \theta \vec{e}_x + \sin \theta \vec{e}_y$ and $\vec{e}_\theta = -r \sin \theta \vec{e}_x + r \cos \theta \vec{e}_y$

Let us differentiate:

$$\frac{\partial \vec{f}(r, \theta)}{\partial \theta} = \frac{\partial \cos \theta}{\partial \theta} \vec{e}_r - \frac{1}{r} \frac{\partial \sin \theta}{\partial \theta} \vec{e}_\theta = -\sin \theta \vec{e}_r - \frac{\cos \theta}{r} \vec{e}_\theta$$

OK, fine!

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OK, fine!

Wait! $\vec{f}(r, \theta) = \cos \theta \vec{e}_r - \frac{\sin \theta}{r} \vec{e}_\theta$

$$= \cos \theta (\cos \theta \vec{e}_x + \sin \theta \vec{e}_y) - \frac{\sin \theta}{r} (-r \sin \theta \vec{e}_x + r \cos \theta \vec{e}_y)$$

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$$= \vec{e}_x = \overrightarrow{\text{constant!}} \text{ How can its derivative be } \neq \vec{0}?$$

Second motivation: another issue with polar coordinates

For instance: $\vec{f}(r, \theta) = \cos \theta \vec{e}_r - \frac{\sin \theta}{r} \vec{e}_\theta$

with (slide 6) $\vec{e}_r = \cos \theta \vec{e}_x + \sin \theta \vec{e}_y$ and $\vec{e}_\theta = -r \sin \theta \vec{e}_x + r \cos \theta \vec{e}_y$

Let us differentiate: are also functions of (r, θ) !

~~$$\frac{\partial \vec{f}(r, \theta)}{\partial \theta} = \frac{\partial \cos \theta}{\partial \theta} \vec{e}_r - \frac{1}{r} \frac{\partial \sin \theta}{\partial \theta} \vec{e}_\theta = -\sin \theta \vec{e}_r - \frac{\cos \theta}{r} \vec{e}_\theta$$~~

Use the product rule!

OK, fine!

Wait! $\vec{f}(r, \theta) = \cos \theta \vec{e}_r - \frac{\sin \theta}{r} \vec{e}_\theta$

$$= \cos \theta (\cos \theta \vec{e}_x + \sin \theta \vec{e}_y) - \frac{\sin \theta}{r} (-r \sin \theta \vec{e}_x + r \cos \theta \vec{e}_y)$$

$$= \vec{e}_x = \overrightarrow{\text{constant!}}$$

How can its derivative be $\neq \vec{0}$

Second motivation: another issue with polar coordinates

For instance: vector field $\vec{c}(r, \theta) = c^r(r, \theta)\vec{e}_r(r, \theta) + c^\theta(r, \theta)\vec{e}_\theta(r, \theta)$

To differentiate w.r.t. r or θ , one must not forget to differentiate the basis vectors...

Every time? But that is time-consuming!

There is a nice trick!

Covariant derivative

Consider the (position dependent: P) basis vectors $\{\vec{e}_i(P)\}$ associated with coordinates $\{x^i\}$.

The partial derivative $\frac{\partial \vec{e}_i(P)}{\partial x^j}$ is itself a vector at P :

☞ can be written as linear combination of the $\{\vec{e}_i(P)\}$:

$$\frac{\partial \vec{e}_i(P)}{\partial x^j} = \Gamma_{ij}^k(P) \vec{e}_k(P)$$

with coefficients Γ_{ij}^k : **Christoffel symbols.** (symmetric under $i \leftrightarrow j$)

Consider now a vector field $\vec{c}(P) = c^i(P) \vec{e}_i(P)$

One has $\frac{\partial \vec{c}(P)}{\partial x^j} = \frac{dc^i(P)}{dx^j} \vec{e}_i(P)$ with

$$\frac{dc^i(P)}{dx^j} \equiv \frac{\partial c^i(P)}{\partial x^j} + \Gamma_{jk}^i(P) c^k(P)$$

Covariant derivative

For a vector field $\vec{c}(P) = c^i(P)\vec{e}_i(P)$ one has $\frac{\partial \vec{c}(P)}{\partial x^j} = \frac{dc^i(P)}{dx^j}\vec{e}_i(P)$ with the covariant derivatives

$$\frac{dc^i(P)}{dx^j} \equiv \frac{\partial c^i(P)}{\partial x^j} + \Gamma_{jk}^i(P)c^k(P)$$

Proof:

$$\frac{\partial \vec{c}(P)}{\partial x^j} = \frac{\partial c^i(P)}{\partial x^j}\vec{e}_i(P) + c^i(P)\frac{\partial \vec{e}_i(P)}{\partial x^j} = \frac{\partial c^i(P)}{\partial x^j}\vec{e}_i(P) + \underbrace{c^i(P)\Gamma_{ij}^k(P)\vec{e}_k(P)}_{c^i\Gamma_{ij}^k\vec{e}_k = c^k\Gamma_{jk}^i\vec{e}_i}$$

$$\frac{\partial \vec{c}(P)}{\partial x^j} = \frac{\partial c^i(P)}{\partial x^j}\vec{e}_i(P) + c^k(P)\Gamma_{jk}^i(P)\vec{e}_i(P) = \frac{dc^i(P)}{dx^j}\vec{e}_i(P)$$

Covariant derivative: polar coordinates

From $\vec{e}_r(r, \theta) = \cos \theta \vec{e}_x + \sin \theta \vec{e}_y$, $\vec{e}_\theta(r, \theta) = -r \sin \theta \vec{e}_x + r \cos \theta \vec{e}_y$

one computes

$$\frac{\partial \vec{e}_r(r, \theta)}{\partial r} = \vec{0} , \quad \frac{\partial \vec{e}_r(r, \theta)}{\partial \theta} = \frac{1}{r} \vec{e}_\theta(r, \theta) , \quad \frac{\partial \vec{e}_\theta(r, \theta)}{\partial r} = \frac{1}{r} \vec{e}_\theta(r, \theta) , \quad \frac{\partial \vec{e}_\theta(r, \theta)}{\partial \theta} = -r \vec{e}_r(r, \theta)$$

These derivatives are linear combinations $\frac{\partial \vec{e}_i(P)}{\partial x^j} = \Gamma_{ij}^k(P) \vec{e}_k(P)$ with the Christoffel symbols

$$\Gamma_{rr}^r = \Gamma_{rr}^\theta = 0, \quad \Gamma_{r\theta}^\theta = \Gamma_{\theta r}^\theta = \frac{1}{r}, \quad \Gamma_{\theta\theta}^r = -r, \quad \Gamma_{r\theta}^r = \Gamma_{\theta r}^r = 0, \quad \Gamma_{\theta\theta}^\theta = 0$$

$$\left\{ \begin{array}{l} \frac{dc^i}{dr} = \frac{\partial c^i}{\partial r} + \Gamma_{rk}^i c^k = \frac{\partial c^i}{\partial r} + \Gamma_{r\theta}^i c^\theta = \frac{\partial c^i}{\partial r} + \frac{c^\theta}{r} \delta^{i\theta} \\ \frac{dc^i}{d\theta} = \frac{\partial c^i}{\partial \theta} + \Gamma_{\theta k}^i c^k = \frac{\partial c^i}{\partial \theta} + \Gamma_{\theta r}^i c^r + \Gamma_{\theta\theta}^i c^\theta = \frac{\partial c^i}{\partial \theta} + \frac{c^r}{r} \delta^{i\theta} - r c^\theta \delta^{ir} \end{array} \right.$$

Covariant derivative: polar coordinates

Consider again the example

$$\vec{c}(r, \theta) = \cos \theta \vec{e}_r - \frac{\sin \theta}{r} \vec{e}_\theta \quad (= \vec{e}_x)$$

i.e. with components $c^r = \cos \theta$, $c^\theta = -\frac{\sin \theta}{r}$

and let us compute the covariant derivatives w.r.t r and θ :

● w.r.t r : $\frac{dc^i}{dr} = \frac{\partial c^i}{\partial r} + \frac{c^\theta}{r} \delta^{i\theta}$ gives $\frac{dc^r}{dr} = \frac{\partial c^r}{\partial r} = 0$, $\frac{dc^\theta}{dr} = \frac{\partial c^\theta}{\partial r} + \frac{c^\theta}{r} = 0$

● w.r.t θ : $\frac{dc^i}{d\theta} = \frac{\partial c^i}{\partial \theta} + \frac{c^r}{r} \delta^{i\theta} - c^\theta \delta^{ir}$ gives

$$\frac{dc^r}{d\theta} = \frac{\partial c^r}{\partial \theta} - r c^\theta = 0 \quad \text{and} \quad \frac{dc^\theta}{d\theta} = \frac{\partial c^\theta}{\partial \theta} + \frac{c^r}{r} = 0$$

All covariant derivatives vanish!



Covariant derivative: rules for arbitrary tensor fields

● Scalar field: $\frac{dc(P)}{dx^j} \equiv \frac{\partial c(P)}{\partial x^j}$

● "Contravariant vector": $\frac{dc^i(P)}{dx^j} \equiv \frac{\partial c^i(P)}{\partial x^j} + \Gamma_{jk}^i(P)c^k(P)$

● "Covariant vector": $\frac{dc_i(P)}{dx^j} \equiv \frac{\partial c_i(P)}{\partial x^j} - \Gamma_{ij}^k(P)c_k(P)$

● Arbitrary (m contra-, n covariant indices) tensor field:

$$\begin{aligned} \frac{d\mathbf{T}_{j_1 \dots j_n}^{i_1 \dots i_m}(P)}{dx^k} &= \frac{\partial \mathbf{T}_{j_1 \dots j_n}^{i_1 \dots i_m}(P)}{\partial x^k} + \Gamma_{kl}^{i_1}(P) \mathbf{T}_{j_1 \dots j_n}^{li_2 \dots i_m}(P) + \dots + \Gamma_{kl}^{i_m}(P) \mathbf{T}_{j_1 \dots j_n}^{i_1 \dots i_{m-1}l}(P) \\ &\quad - \Gamma_{j_1 k}^l(P) \mathbf{T}_{lj_2 \dots j_n}^{i_1 \dots i_m}(P) - \dots - \Gamma_{j_n k}^l(P) \mathbf{T}_{j_1 \dots j_{n-1}l}^{i_1 \dots i_m}(P) \end{aligned}$$

A last formula

To compute the Christoffel symbols with the help of the metric tensor and its inverse 🖱️ no need to go back to the expression of the basis vectors $\{\vec{e}_i\}$ in terms of those of the Cartesian / Minkowski basis.

$$\Gamma_{ij}^k(P) = \frac{1}{2} g^{kl}(P) \left[\frac{\partial g_{jl}(P)}{\partial x^i} + \frac{\partial g_{il}(P)}{\partial x^j} - \frac{\partial g_{ij}(P)}{\partial x^l} \right]$$

🖱️ easy to automatize (symbolic algebra programs)