

Tutorial sheet 13

The exercise marked with a star is homework.

*34. Speed of sound in ultrarelativistic matter

Consider a perfect fluid with the usual energy-momentum tensor. $T^{\mu\nu} = \mathcal{P}g^{\mu\nu} + (\epsilon + \mathcal{P})u^\mu u^\nu / c^2$. It is assumed that there is no conserved quantum number relevant for thermodynamics, so that the energy density in the local rest frame ϵ is function of a single thermodynamic variable, for instance $\epsilon = \epsilon(\mathcal{P})$. Throughout the exercise, Minkowski coordinates are used.

A background “flow” with uniform local-rest-frame energy density and pressure ϵ_0 and \mathcal{P}_0 is submitted to a small perturbation resulting in $\epsilon = \epsilon_0 + \delta\epsilon$, $\mathcal{P} = \mathcal{P}_0 + \delta\mathcal{P}$, and $\vec{v} = \vec{0} + \delta\vec{v}$.

i. Starting from the energy-momentum conservation equation $\partial_\mu T^{\mu\nu} = 0$, show that linearization to first order in the perturbations leads to the two equations of motion $\partial_t \delta\epsilon = -(\epsilon_0 + \mathcal{P}_0)\vec{\nabla} \cdot \delta\vec{v}$ and $(\epsilon_0 + \mathcal{P}_0)\partial_t \delta\vec{v} = -c^2 \vec{\nabla} \delta\mathcal{P}$.

ii. Show that the speed of sound is given by the expression $c_s^2 = \frac{c^2}{d\epsilon/d\mathcal{P}}$.

iii. Compute c_s for a fluid obeying the Stefan–Boltzmann law¹ $\mathcal{P} = \frac{g\pi^2}{90} \frac{(k_B T)^4}{(\hbar c)^3}$, with g the number of degrees of freedom (e.g. $g = 2$ for blackbody radiation).

Hint: You may find the Gibbs–Duhem relation useful...

¹This is a good opportunity to refresh your knowledge on the statistical physics of relativistic systems. Can you give a physical argument why quantum effects always play a role in such systems, as signaled by the presence of \hbar in the equation of state?