Tutorial sheet 4

Throughout this exercise sheet we use Minkowski coordinates $\{x^{\mu}\}.$

7. Non-relativistic limit

Let $u^{\mu}(x)$ denote the four-velocity of a fluid. With its help, one defines the four-vector field

$$
a^{\mu}(\mathsf{x}) \equiv u^{\nu}(\mathsf{x}) \partial_{\nu} u^{\mu}(\mathsf{x}). \tag{1}
$$

How can one interpret a^{μ} ? Compute the non-relativistic limit $|\vec{v}| \ll 1 (= c!)$ of a^{μ} , where $v^{i} = u^{i}/u^{0}$ is the 3-velocity associated with u^{μ} . What do you recognize? Is this consistent with your interpretation of a^{μ} ?

8. A family of solutions of the dynamical equations for perfect relativistic fluids

Let $\tau^2 \equiv -x^\mu x_\mu$, where the "mostly plus" metric is used. Show that the following four-velocity, pressure and charge density constitute a solution of the equations describing the motion of a perfect relativistic fluid with equation of state $P = K \varepsilon$ — where K is a constant — and a single conserved charge:

$$
u^{\mu}(\mathsf{x}) = \frac{x^{\mu}}{\tau} \quad , \quad \mathcal{P}(\mathsf{x}) = \mathcal{P}_0 \left(\frac{\tau_0}{\tau}\right)^{3(1+K)} \quad , \quad n(\mathsf{x}) = n_0 \left(\frac{\tau_0}{\tau}\right)^3 \mathcal{N}\big(\sigma(\mathsf{x})\big), \tag{2}
$$

with τ_0 , \mathcal{P}_0 , \mathcal{n}_0 arbitrary constants and N an arbitrary function of a single argument, while σ is a function of spacetime coordinates with vanishing comoving derivative: $u^{\mu}(x)\partial_{\mu}\sigma(x) = 0$.

Discuss the physical meaning of the solution (you may already know the limiting case with no conserved charge, $n = 0$).