

## 7.7 The anomalous magnetic moment of the electron/muon

effective Lagrangian for non-relativistic (NR) electrons and soft ( $q^2 \ll m^2$ ) photons:

$$L_{\text{eff}} = \bar{\chi} \left( i \not{D}_0 + \frac{\vec{D}^2}{2m} - e g \vec{B} \cdot \frac{\vec{\sigma}}{2} \right) \chi \quad \begin{array}{l} \text{non-relativistic QED} \\ \text{(NRQED)} \end{array}$$

Ignoring the effect of hard ( $l^2 \sim m^2$ ) virtual particles we found  $g = 2$

Now we'll compute the effect of hard virtual particles by computing the amplitude for a NR electron scattering off a soft magnetic field in QED and NRQED and adjust  $g$  such that the results match.

$$\left( \text{---} \right)_{\text{NRQED}} = \left( \text{---} + \text{---} \right)_{\text{QED}}$$

To determine  $g$ , we only need the spin-dependent part.

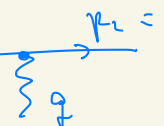
### The amplitude in QED

We only need to include the effect of hard loop momenta  $l$ .

Since we consider  $q^2 \ll l^2$ , we may put  $q^2 = 0$  in  $\bar{u} \sigma \Gamma u$ .

Since  $F_1(0) = 0$ , the only correction comes from  $F_2(0)$ .

$$u := \begin{pmatrix} 1 \\ \vec{0} \end{pmatrix} \quad \vec{q}, \vec{k} = \mathcal{O}(v), \quad q^0, k^0 = \mathcal{O}(v^2)$$

$$p_1 = mu + k \quad p_2 = mu + k + q \quad = \bar{u}(-ie)u$$


$$u(p_1) = \sqrt{2m} \begin{pmatrix} \xi \\ \lambda \end{pmatrix} \quad \lambda = \mathcal{O}(v)$$

$$\not{p}_1 - m = m(\not{t} - 1) + \not{k} = \begin{pmatrix} 0 & -\vec{k} \cdot \vec{\sigma} \\ \vec{k} \cdot \vec{\sigma} & -2m \end{pmatrix} + \mathcal{O}(v^2)$$

$$(\not{p}_1 - m)u = 0 \Rightarrow \vec{k} \cdot \vec{\sigma} - 2m\lambda = 0,$$

$$u_S(p_1) = \sqrt{2m} \begin{pmatrix} \xi_S \\ \frac{\vec{k} \cdot \vec{\sigma}}{2m} \xi_S \end{pmatrix}$$

Since we consider the scattering off a magnetic field, we need the spatial components:

$$\bar{u}_{s_2}(\vec{p}_2)(-ie\gamma^\mu)u_{s_1}(\vec{p}_1) = -ie \underbrace{u^\dagger \gamma^0 \gamma^\mu u}_{= \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & \sigma^\mu \\ -\sigma^\mu & 0 \end{pmatrix} = \begin{pmatrix} 0 & \sigma^\mu \\ \sigma^\mu & 0 \end{pmatrix}}$$

$$= -ie 2m \left( \xi_{s_2}^\dagger, \xi_{s_2}^\dagger \frac{(\vec{k} + \vec{q}) \cdot \vec{\sigma}}{2m} \right) \begin{pmatrix} \sigma^\mu \frac{\vec{k} \cdot \vec{\sigma}}{2m} \xi_{s_1} \\ \sigma^\mu \xi_{s_1} \end{pmatrix}$$

$$= -ie \cdot \xi_{s_2}^\dagger \left[ \sigma^\mu \vec{k} \cdot \vec{\sigma} + (\vec{k} + \vec{q}) \cdot \vec{\sigma} \sigma^\mu \right] \xi_{s_1}$$

$$= \underbrace{\{\sigma^\mu, \vec{k} \cdot \vec{\sigma}\}}_{\text{spin-indep.}} + \vec{q} \cdot \vec{\sigma} \sigma^\mu = 2k^\mu + \frac{1}{2} \{\vec{q} \cdot \vec{\sigma}, \sigma^\mu\} + \frac{1}{2} [\vec{q} \cdot \vec{\sigma}, \sigma^\mu]$$

$$= \underbrace{2k^\mu + q^\mu}_{\text{spin-indep.}} + \underbrace{i\epsilon^{jmn} q^j \sigma^n}_{\text{spin dependent}}$$

$$\left( \begin{array}{c} \text{---} \bullet \text{---} \\ | \\ \text{---} \end{array} \right)_{\text{spin dependent}} = e \epsilon^{\mu\nu j} q^\nu \tilde{S}_{\lambda_2}^\dagger \sigma^j \tilde{S}_{\lambda_1}$$

The contribution from  $F_2(0)$ :

$$\frac{\text{---} \bullet \text{---}}{\text{---}} = -ie \bar{u} F_2 \frac{i\sigma^{\mu\nu} q^\nu}{2m} u \approx \frac{eF_2}{2m} (-1) q^\nu \bar{u} \sigma^{\mu\nu} u$$

$$\begin{aligned} \sigma^{\mu\nu} &= \frac{i}{2} \begin{pmatrix} 0 & \sigma^{\mu\nu} \\ -\sigma^{\mu\nu} & 0 \end{pmatrix} - (\mu \leftrightarrow \nu) = \frac{i}{2} \begin{pmatrix} -[\sigma^\mu, \sigma^\nu] & 0 \\ 0 & -[\sigma^\mu, \sigma^\nu] \end{pmatrix} \\ &= \epsilon^{\mu\nu j} \begin{pmatrix} \sigma^j & 0 \\ 0 & \sigma^j \end{pmatrix} \end{aligned}$$

$$\bar{u} \sigma^{\mu\nu} u = \epsilon^{\mu\nu j} 2m \tilde{S}_{\lambda_2}^\dagger \sigma^j \tilde{S}_{\lambda_1} \Rightarrow$$

$$\left( \begin{array}{c} \text{---} \bullet \text{---} \\ | \\ \text{---} \end{array} + \frac{\text{---} \bullet \text{---}}{\text{---}} \right)_{\text{spin dependent}} = e \epsilon^{\mu\nu j} q^\nu \tilde{S}_{\lambda_2}^\dagger \sigma^j \tilde{S}_{\lambda_1} [1 + F_2(0)]$$

The 1 in  $(1 + F_2)$  is reproduced in NRQED by  $g=2 \Rightarrow$

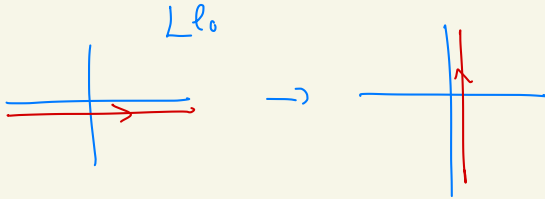
$$g = 2 [1 + F_2(0)]$$

compute  $F_2(0)$

$$F_2(0) = i \delta e^2 m^2 \int_0^1 dx \int_0^1 dy \int_0^1 dz \delta(1-x-y-z) z(1-z) \int_{\mathcal{D}} \frac{1}{\mathcal{D}^3}$$

where now

$$\mathcal{D} = l^2 - M^2 \quad \text{with} \quad M = (1-z)m$$

Wick rotation: 

$$\int_{\mathcal{D}} \frac{1}{(l^2 - M^2)^3} = i (-1)^3 \int_{\mathcal{D}} \frac{1}{(l^2 + M^2)^3} \stackrel{\text{sec. 4.4}}{=} -i \frac{M^{-2}}{(4\pi)^2} \frac{\Gamma(3-2)}{\Gamma(3)}$$

$$= -\frac{i}{(4\pi)^2 M^2} \frac{1}{2 \cdot 1 \cdot \Gamma(1)} = \frac{-i}{(4\pi)^2} \frac{1}{2M^2}$$

$$F_2(0) = i \delta e^2 \frac{-i}{2(4\pi)^2} \int_0^1 dx \int_0^1 dy \int_0^1 dz \delta(1-x-y-z) \frac{z}{1-z}$$

$$= \frac{\alpha}{\pi} \int_0^1 dz \frac{z}{1-z} \int_0^{1-z} dy = \frac{\alpha}{2\pi} \quad \Rightarrow$$

$$\boxed{g_{-2} = \frac{\alpha}{\pi}}$$